Electrical resistivity in comparison to thermal conductivity of selected alloys

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Motivation and task

- Every month, several tons of spacecraft mass enter uncontrolledly the atmosphere
- International agreements require for each satellite launched into LEO to conduct either a controlled de-orbiting or to assess the possible risk for human population
- The demisability of a spacecraft and its components during re-entry depends on many parameters
- Numerical simulation of demise at re-entry of a spacecraft requires among others thermophysical properties up to complete melting of decay
- Common alloys for aerospace applications: TiAl6V4, Aluminium, stainless steel 316L, superalloys...
- A full set of thermophysical properties in the range room temperature to melting was measured to characterise the alloy:
 - Electrical resistivity
 - Specific heat capacity
 - Density and linear thermal expansion
 - Thermal diffusivity thermal conductivity
 - A comparison between thermal and electrical conductivity/resistivity was made to check the validity of the Wiedemann-Franz law

Numerical simulation of demise at re-entry of Gensat



Measurement of heat capacity

Measurement:

- **NETZSCH DSC 404 Pegasus heat-flux** differential-scanning calorimeter
- Empty pan, NIST SRM 781 or sapphire, alloy
- Specimens of approx. 100 to 300 mg
- Heating and cooling rates of 10 K/min
- Platinum crucibles (partly Alumina-liner)



Measurement of thermal diffusivity



Measurement:

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- NETZSCH LFA 427
- Specimens: 12.5 mm in diameter;
 3 mm thickness
- Sand-blasted, UHV
- 500 μs laser pulse length
- Cape-Lehmann fit with pulse length and heat loss correction

Measurement of thermal expansion and density



- NETZSCH DIL 402 CD
- Specimens: 6 mm in diameter; 25mm in length
- Heating/cooling rate 2 K/min
- Density at room temperature by an Archimedean balance (Sartorius ED224S)

Calculation of density and thermal conductivity

• Density ρ as a function of temperature T is calculated by density at room temperature ρ_0 and thermal expansion $\Delta l/l_0$:

$$\rho(T) = \frac{\rho_0}{(1 + \frac{\Delta l(T)}{l_0})^3}$$

• Thermal conductivity λ as a function of temperature T is calculated by thermal diffusivity a_0 , density ρ , and heat capacity c_P :

$$\lambda(T) = \left[a_0(T) \left(1 + \frac{\Delta l(T)}{l_0}\right)^2\right] \rho(T) c_P(T)$$

- Millisecond pulse heating system capable to provide up to 5000 A of current to heat a metallic specimen
- Four-probe measurement of current and voltage along a specific portion of the specimen
- Self-heating of the specimen, (almost) all imparted electrical energy is used to increase enthalpy of the specimen (long, thin rod approximation)
- > But:

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- Specimen has a certain thickness (4 mm diameter)
- At elevated temperatures, there are radiation losses
- Steel low thermal conductivity (radiation), aluminium high thermal conductivity (conduction)
- Long, thin rod approximation is violated
- Deviation from the long, thin rod approximation by heat loss is computed using a numerical simulation model
- From the heat capacity measurements, the enthalpy vs. temperature relation is known (no use of pyrometer no emissivity problem)

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1 interferometer, 2 vessel, 3 pyrometer





Calculation of electrical resistivity and enthalpy

Electrical resistivity ρ_{el} as a function of time t is calculated by voltage U, current I, diameter d, length between the knife edges l, and thermal expansion $\Delta l/l_0$

$$\rho_{el}(t) = \frac{U(t) d^2 \pi}{4 l(t) l} (1 + \frac{\Delta l}{l_0})$$

Specific enthalpy H as a function of time t is calculated by thermal voltage U, current I, and mass between the knife edges m :

$$H(t) = \frac{1}{m} \int_0^t U(\tau) I(\tau) d\tau$$





Temperature is calculated from results of the DSC measurement

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Very good reproducibility between different specimen





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Worst case scenario

Calculation thermal conductivity - electrical resistivity

For the conductivity λ as a function of temperature T is calculated from electrical resistivity ρ_{el} ; the Sommerfeld-value of the Lorenz number L_0 (2.445 x 10⁻⁸ V²/K²), the phonon contribution of thermal conductivity λ_G , and A, a factor to consider scattering at solute atoms and :

Wiedemann – Franz:
$$\lambda(T) = \frac{L_0 T}{\rho_{el}}$$

Smith – Palmer:
$$\lambda(T) = \frac{A L_0 T}{\rho_{el}} + \lambda_G$$

Smith-Palmer-Plot

For the restrict the two probability λ as a function of temperature T is calculated from electrical resistivity ρ_{el} ; the Sommerfeld-value of the Lorenz number L_0 (2.445 x 10⁻⁸ V²/K²), the phonon contribution of thermal conductivity λ_G , and A, a factor to consider scattering at solute atoms and :



Smith-Palmer-Plot

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Smith-Palmer-Plot: Inconel 625



Smith-Palmer-Plot: Inconel 625



Thermal conductivity as a function of temperature: Inconel 625



Smith-Palmer-Plot (forged Ti-6Al-4V):



Results of thermal conductivity (both types of Ti-6AI-4V):



Thermal conductivity as a function of temperature: AISi7Mg - annealed



Thermal conductivity as a function of temperature: AISi7Mg – as cast



Conclusions

- A full set of thermophysical properties in the range room temperature to melting was measured to characterise aerospace alloys:
 - Electrical resistivity
 - Specific heat capacity
 - Density and linear thermal expansion
 - Thermal diffusivity thermal conductivity
- Electrical resistivity can be measured precisely by millisecond pulse heating, heat losses are negligible
- Only a modified Wiedemann-Franz law can be used to calculate thermal conductivity from electrical resistivity - especially when the lattice contribution is relatively high

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